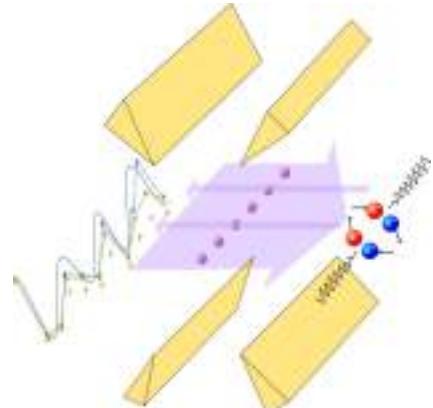


Digital Quantum Simulation and Symmetry Protection with Trapped Ions

Norbert M. Linke

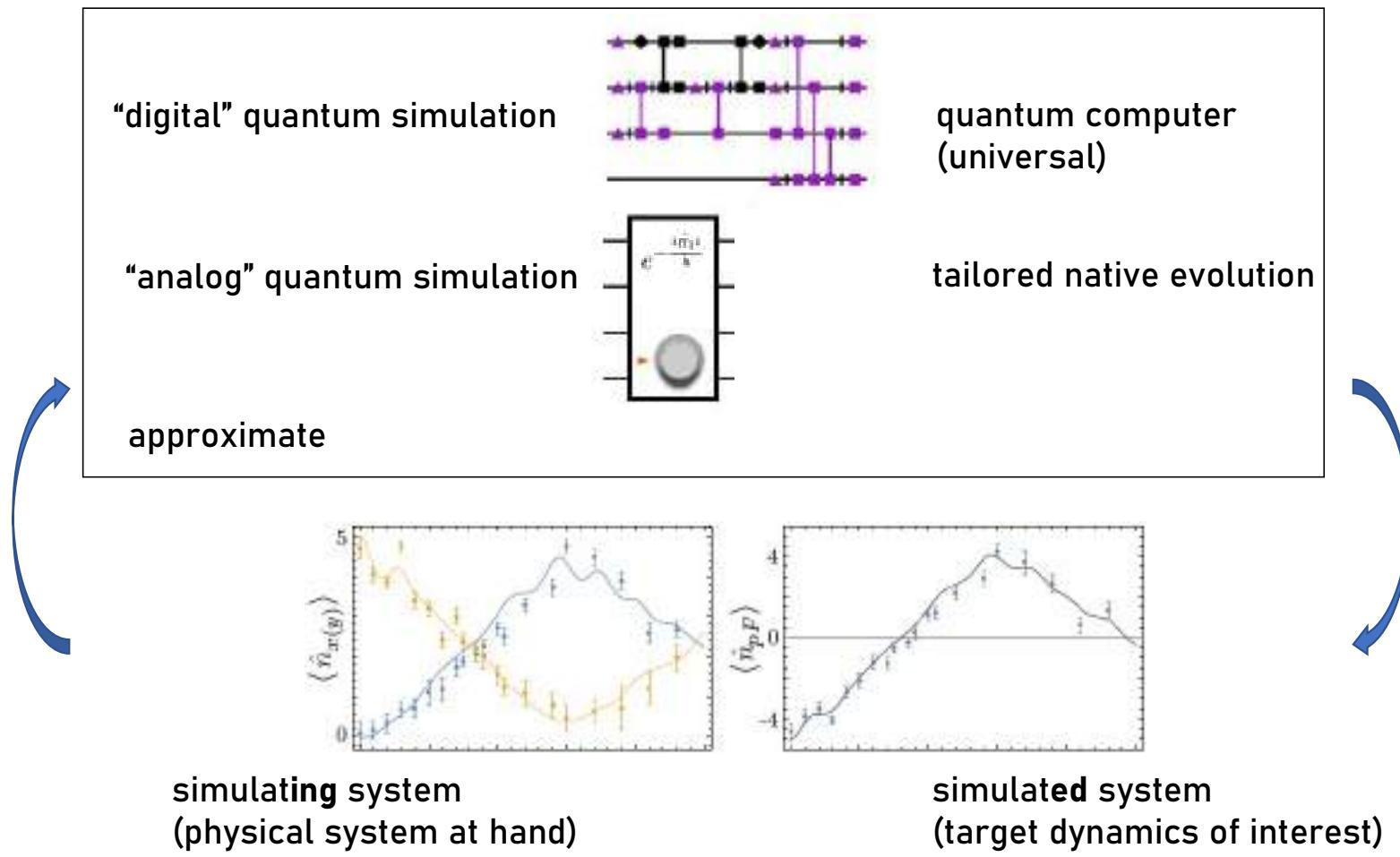
Duke Quantum Center, Duke University, Durham, North Carolina, USA
Joint Quantum Institute, University of Maryland, College Park, Maryland, USA

ECTI, Schloss Bückeburg, Germany
September 25, 2023



Duke Quantum Center

Quantum Simulation

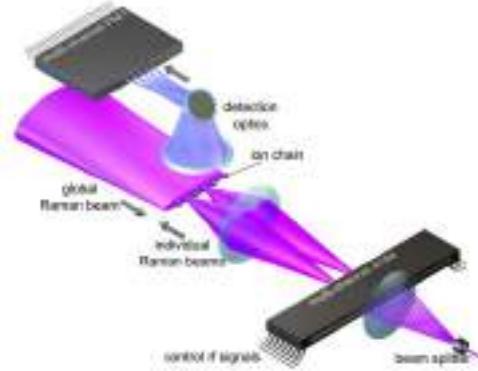


Overview

Experimental system

Individually-addressed $^{171}\text{Yb}^+$ ions

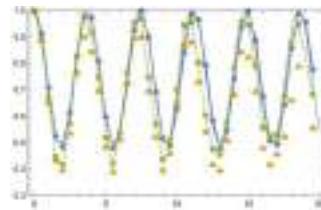
Modular operations and compiler (5-9 qubits)



Applications

Digital Quantum Simulation

The Schwinger model



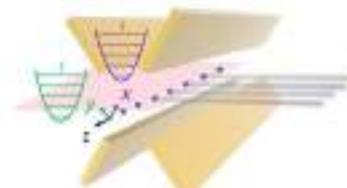
N. H. Nguyen et al., PRX Quantum 3, 020324 (2022).

Hybrid Quantum Simulation

The Yukawa model

Pairwise-parallel gates

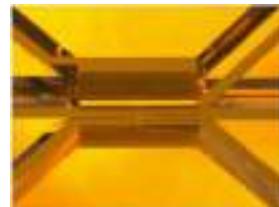
Orthogonal modes



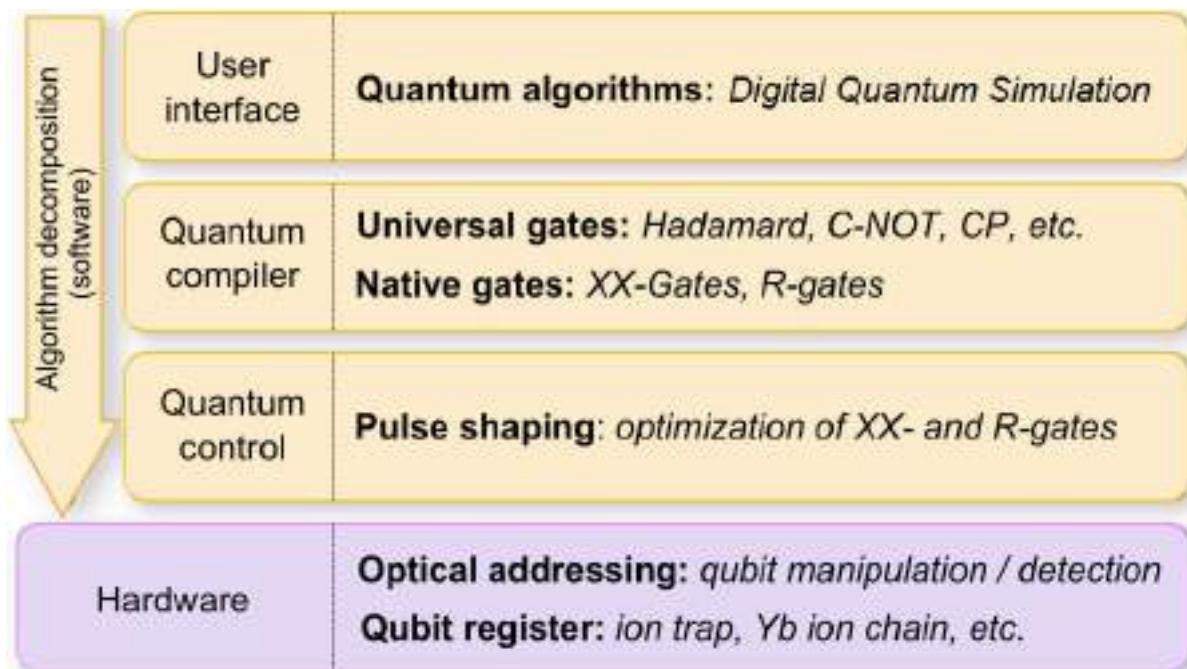
Y. Zhu et al., Adv. Quantum Technol. 020324 (2023).

Outlook

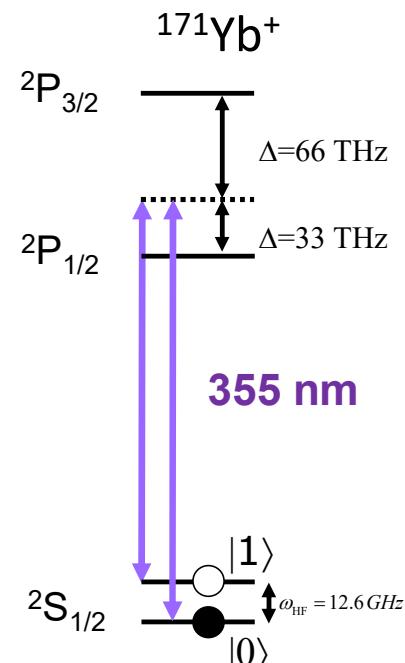
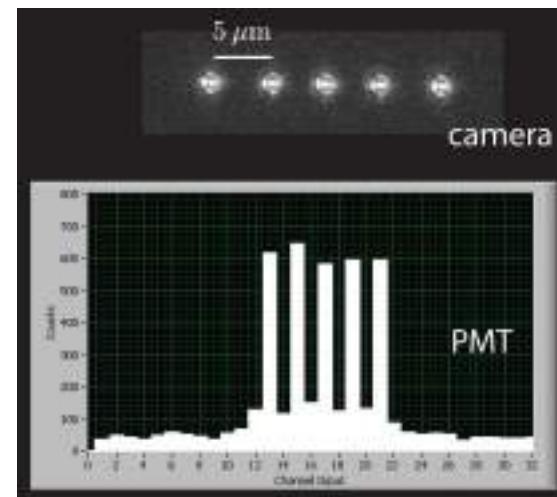
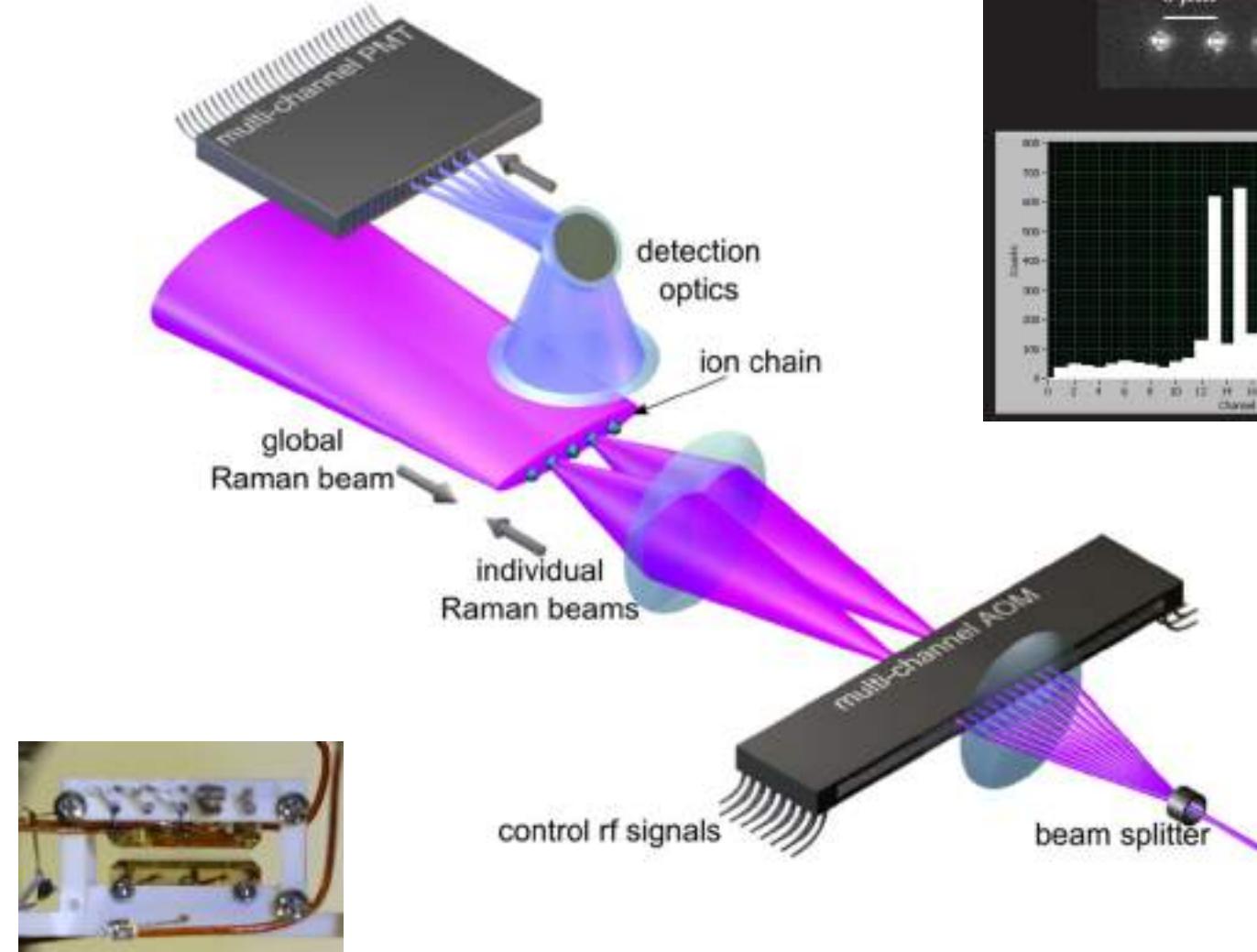
A new monolithic ion trap



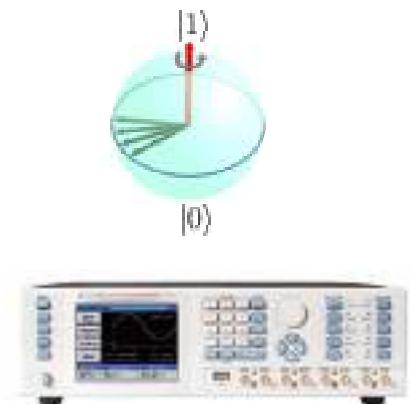
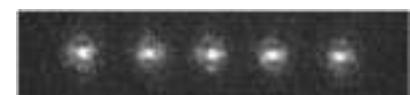
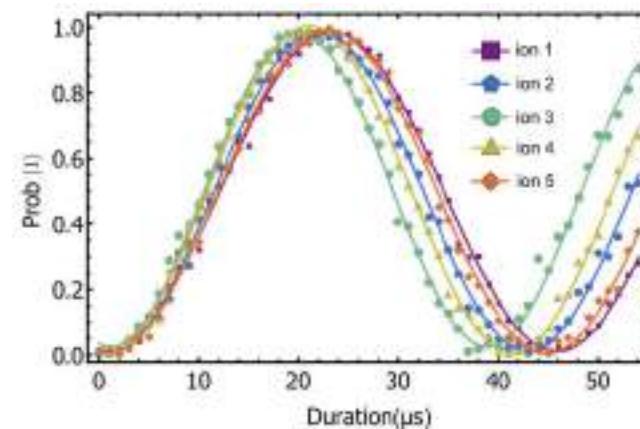
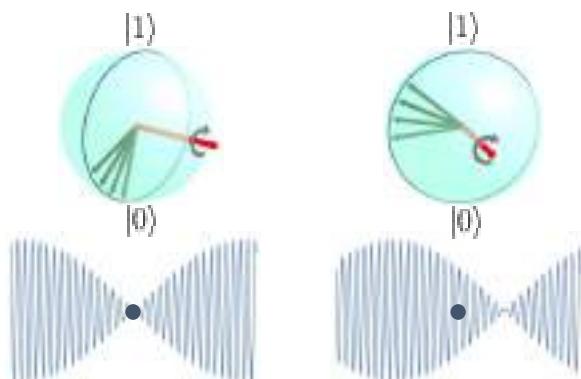
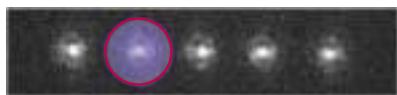
Experimental system: QC architecture



Experimental system: Hardware



Experimental system: Single qubit gates



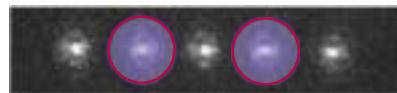
R-gate (x/y rotations)

$$R_\phi(\theta) = \begin{bmatrix} \cos(\frac{\theta}{2}) & -i\sin(\frac{\theta}{2})e^{-i\phi} \\ -i\sin(\frac{\theta}{2})e^{i\phi} & \cos(\frac{\theta}{2}) \end{bmatrix}$$

z-gates

$$R_z(\theta) = \begin{bmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{bmatrix}$$

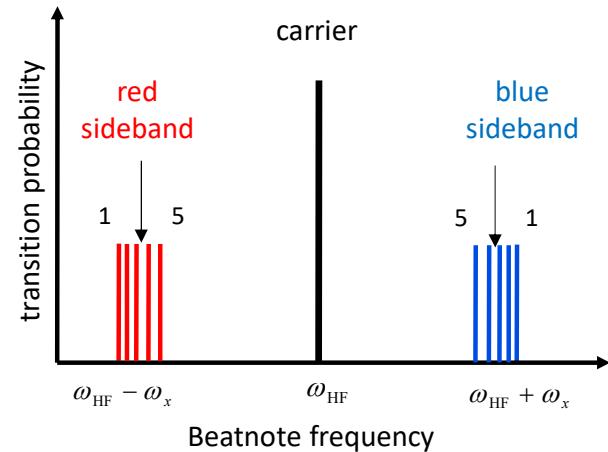
Experimental system: Two-qubit gates



$$U(t) = \exp[-i \sum_{n,k} \hat{D}(\alpha_n^k(t)) \sigma_x^n - i \sum_{i,j} \chi_{ij}(t) \sigma_x^i \sigma_x^j]$$

$$XX(\chi_{i,j}) = \begin{bmatrix} \cos(\chi_{i,j}) & 0 & 0 & -i\sin(\chi_{i,j}) \\ 0 & \cos(\chi_{i,j}) & -i\sin(\chi_{i,j}) & 0 \\ 0 & -i\sin(\chi_{i,j}) & \cos(\chi_{i,j}) & 0 \\ -i\sin(\chi_{i,j}) & 0 & 0 & \cos(\chi_{i,j}) \end{bmatrix}$$

$$|00\rangle \rightarrow \frac{1}{\sqrt{2}}(|00\rangle - i|11\rangle)$$



MS gate: PRL 82 (1999)

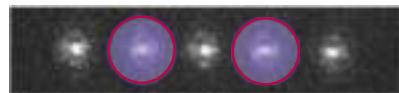
T. Choi et al. PRL 112, 19502 (2014)

T. J. Green et al., PRL 114, 120502 (2015)

P. H. Leung et al. PRL 120, 020501 (2018)

Y. Shapira et al., PRL 121, 180502 (2018)

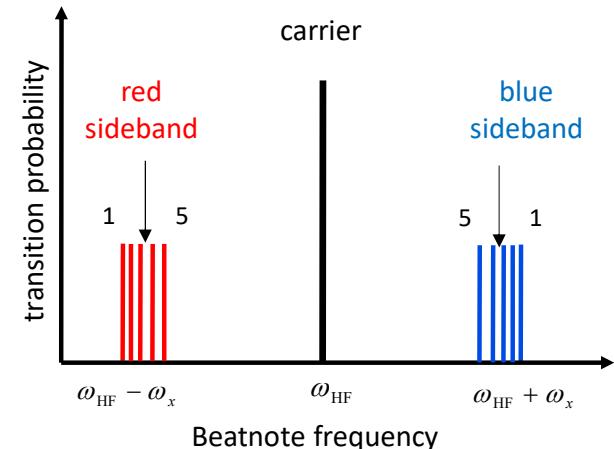
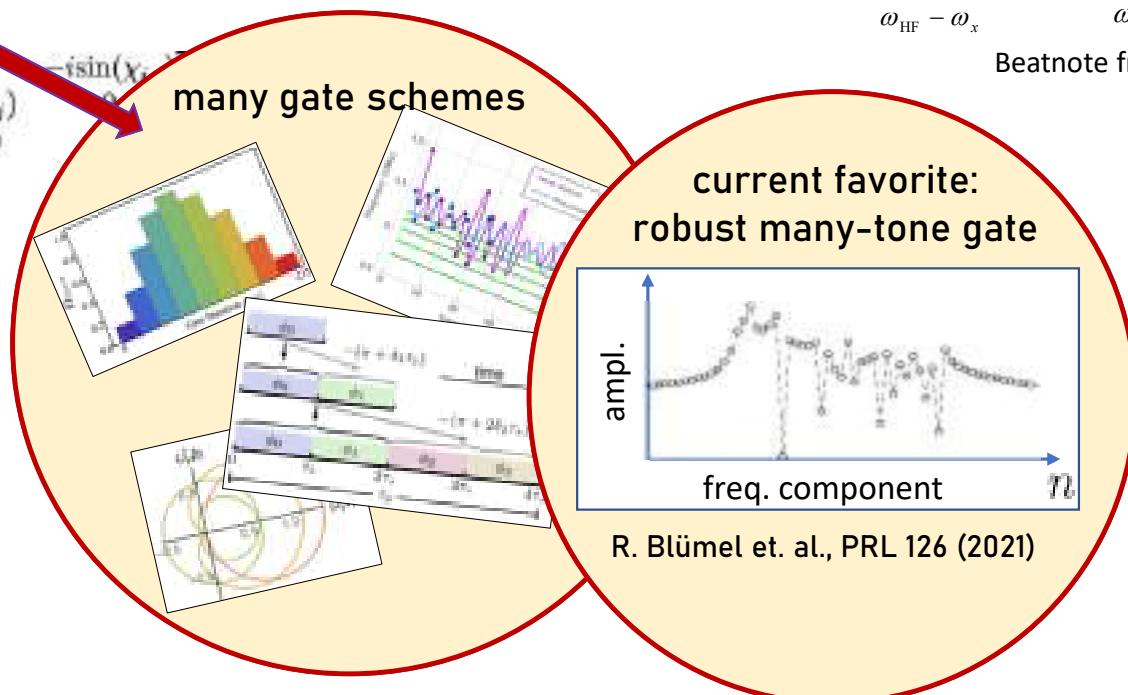
Two-qubit gates



$$U(t) = \exp[-i \sum_{n,k} \hat{D}(\alpha_n^k(t)) \sigma_x^n - i \sum_{i,j} \chi_{ij}(t) \sigma_x^i \sigma_x^j]$$

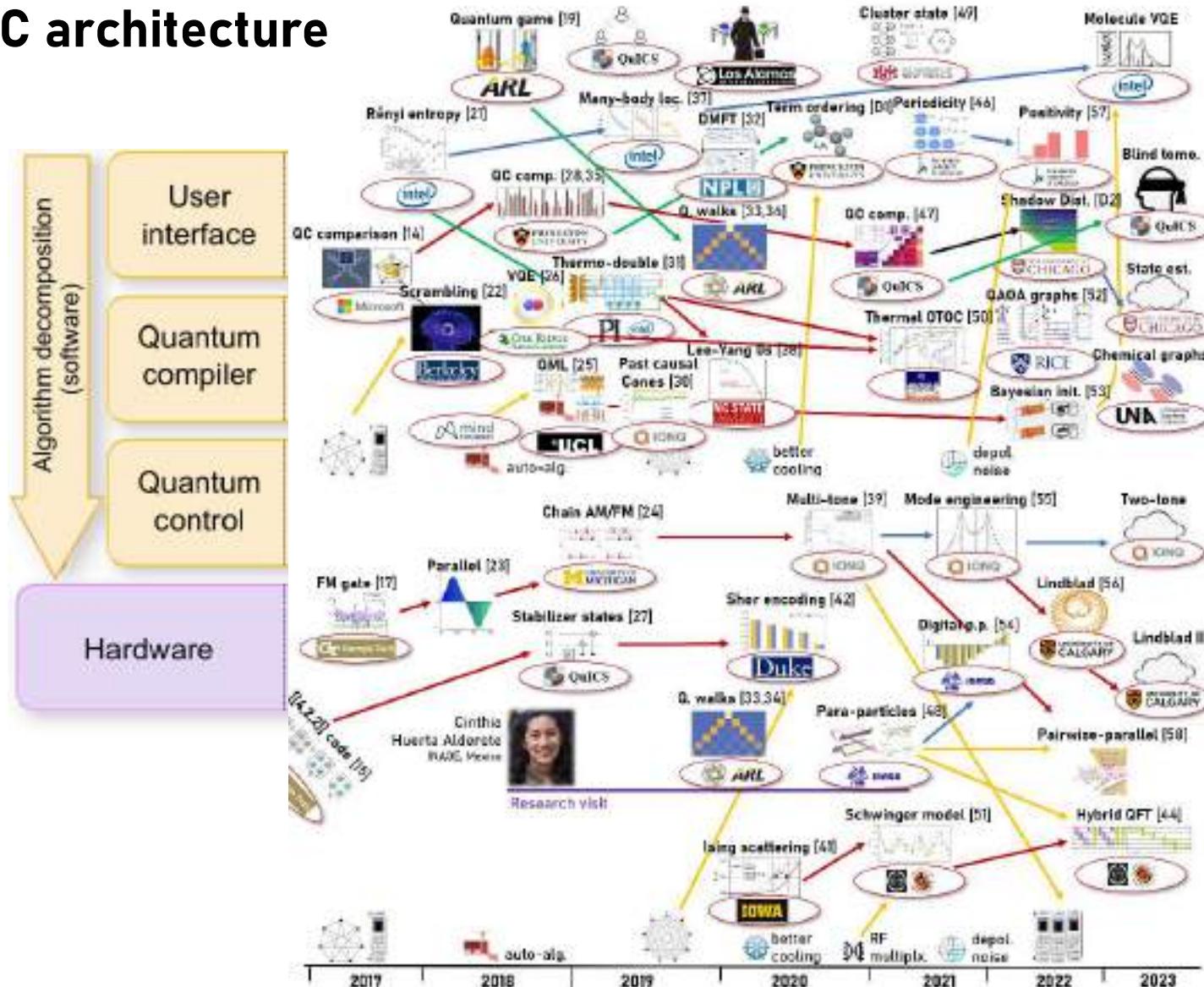
$$XX(\chi_{i,j}) = \begin{bmatrix} \cos(\chi_{i,j}) & 0 & 0 \\ 0 & \cos(\chi_{i,j}) & -i\sin(\chi_{i,j}) \\ 0 & -i\sin(\chi_{i,j}) & \cos(\chi_{i,j}) \end{bmatrix}$$

$$|00\rangle \rightarrow \frac{1}{\sqrt{2}}(|00\rangle - i|11\rangle)$$



MS gate: PRL 82 (1999)
 T. Choi et al. PRL 112, 19502 (2014)
 T. J. Green et al., PRL 114, 120502 (2015)
 P. H. Leung et al. PRL 120, 020501 (2018)
 Y. Shapira et al., PRL 121, 180502 (2018)

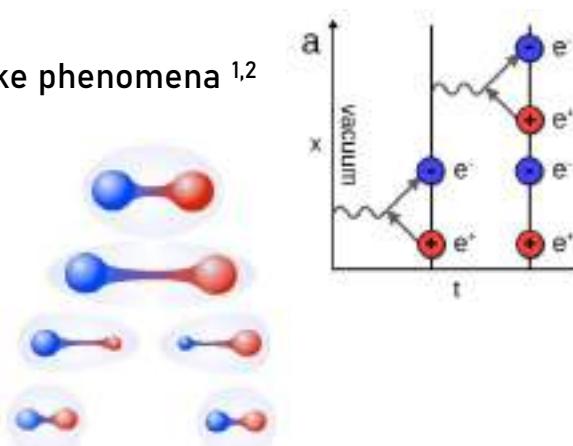
Experimental system: QC architecture



Digital Quantum Simulation: The Schwinger model

Digital Quantum Simulation: The Schwinger model

- Schwinger model is a Quantum Field Theory in 1+1D
- Schwinger model has Quantum Chromodynamics -like phenomena^{1,2}
 - Pair creation-annihilation
 - String breaking
- Testbed for quantum simulation methods^{3,4,5}
- Digital simulation with long(-ish) time dynamics

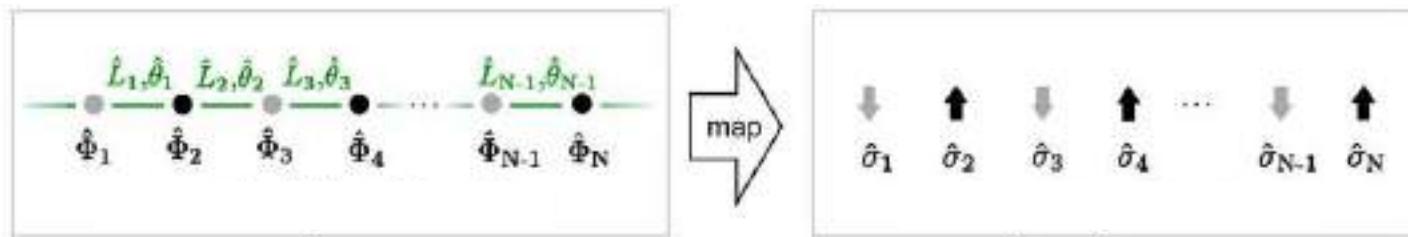


- [1] Coleman Ann. Phys. 101 (1976)
- [2] Hebenstreit et al PRL. 111 (2013)
- [3] Martinez et al Nature 534 (2016)
- [4] Surace et al PRX 10 (2020)
- [5] Mil et al Science 367 (2020)

Digital Quantum Simulation: The Schwinger model

Lattice Schwinger model (spinless 1+1D QFT, discretized space, normalize)

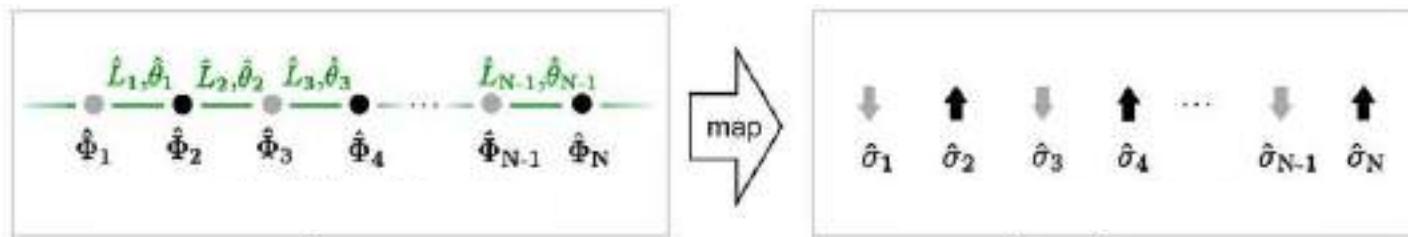
$$\hat{H}_{\text{lat}} = -iw \sum_{n=1}^{N-1} [\hat{\Phi}_n^\dagger e^{i\hat{\theta}_n} \hat{\Phi}_{n+1} - \text{H.C.}] + m \sum_{n=1}^N (-1)^n \hat{\Phi}_n^\dagger \hat{\Phi}_n + J \sum_{n=1}^{N-1} \hat{L}_n^2,$$



Digital Quantum Simulation: The Schwinger model

Lattice Schwinger model (spinless 1+1D QFT, discretized space, normalize)

$$\hat{H}_{\text{lat}} = -iw \sum_{n=1}^{N-1} [\hat{\Phi}_n^\dagger e^{i\hat{\theta}_n} \hat{\Phi}_{n+1} - \text{H.C.}] + m \sum_{n=1}^N (-1)^n \hat{\Phi}_n^\dagger \hat{\Phi}_n + J \sum_{n=1}^{N-1} \hat{L}_n^2,$$



Odd lattice sites:

$$\bullet \cong \text{vac} \cong \uparrow \quad \hat{L}_n = \hat{L}_{n-1}$$

$$\bullet \cong e^+ \cong \downarrow \quad \hat{L}_n = \hat{L}_{n-1} - 1$$

Even lattice sites:

$$\bullet \cong e^- \cong \uparrow \quad \hat{L}_n = \hat{L}_{n-1} + 1$$

$$\bullet \cong \text{vac} \cong \downarrow \quad \hat{L}_n = \hat{L}_{n-1}$$

Gauss' law applies for photon link:

-> number of spin-up qubits conserved

Digital Quantum Simulation: The Schwinger model

Final qubit Hamiltonian

Fermion mass, hopping on lattice, E-field interaction

$$\hat{H}_s = \frac{\mu}{2} \sum_{n=1}^N (-1)^n \sigma_n^z + x \sum_{n=1}^{N-1} \underbrace{\{\sigma_n^+ \sigma_{n+1}^- + \text{h.c.}\}}_{\sigma_n^x \sigma_{n+1}^x + \sigma_n^y \sigma_{n+1}^y} + \frac{1}{4} \sum_{n=1}^{N-1} \left\{ \sum_{m=1}^n \left[\sigma_m^z + (-1)^m \right] \right\}^2$$

parameters

$$x = 0.6, \quad \mu = 0.1, \quad \delta t = 0.5, \quad N_t = 20.$$

Run Hamiltonian evolution as a quantum circuit: Trotterization

$$\hat{U} = e^{-i\hat{H}t/\hbar} = \lim_{n \rightarrow \infty} (\Pi_k e^{-i\hat{H}_k t/n\hbar})^n$$

First-order Trotter approximation: pick finite n

$$\delta t = t/n \quad \hat{U}_k = e^{-i\hat{H}_k \delta t / \hbar} \quad (\text{steps don't commute} \rightarrow \text{term ordering})$$

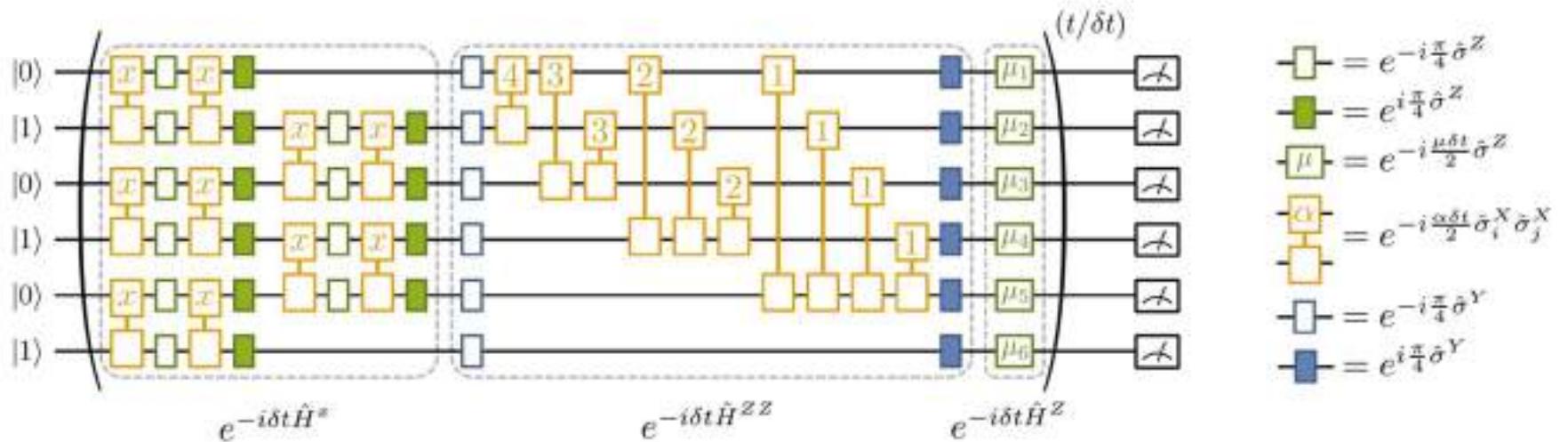
Digital Quantum Simulation: The Schwinger model

Final qubit Hamiltonian

Fermion mass, hopping on lattice, E-field interaction

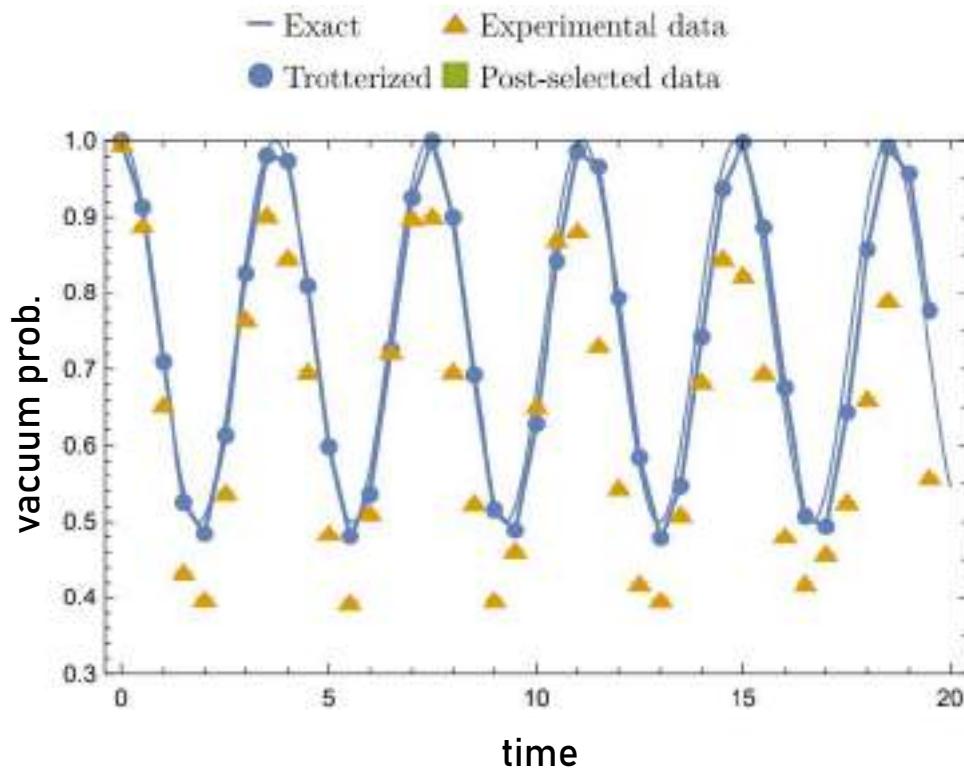
$$\hat{H}_s = \frac{\mu}{2} \sum_{n=1}^N (-1)^n \sigma_n^z + x \sum_{n=1}^{N-1} \{ \sigma_n^+ \sigma_{n+1}^- + \text{h.c.} \} + \frac{1}{4} \sum_{n=1}^{N-1} \left\{ \sum_{m=1}^n \left[\sigma_m^z + (-1)^m \right] \right\}^2$$

- Start at ground state of $x=0$
- Evolve with $x \neq 0$ to time t (Trotterized) $|\text{vac}\rangle = |0101\dots01\rangle$
- Measure vacuum survival prob. / particle number density / E-field density



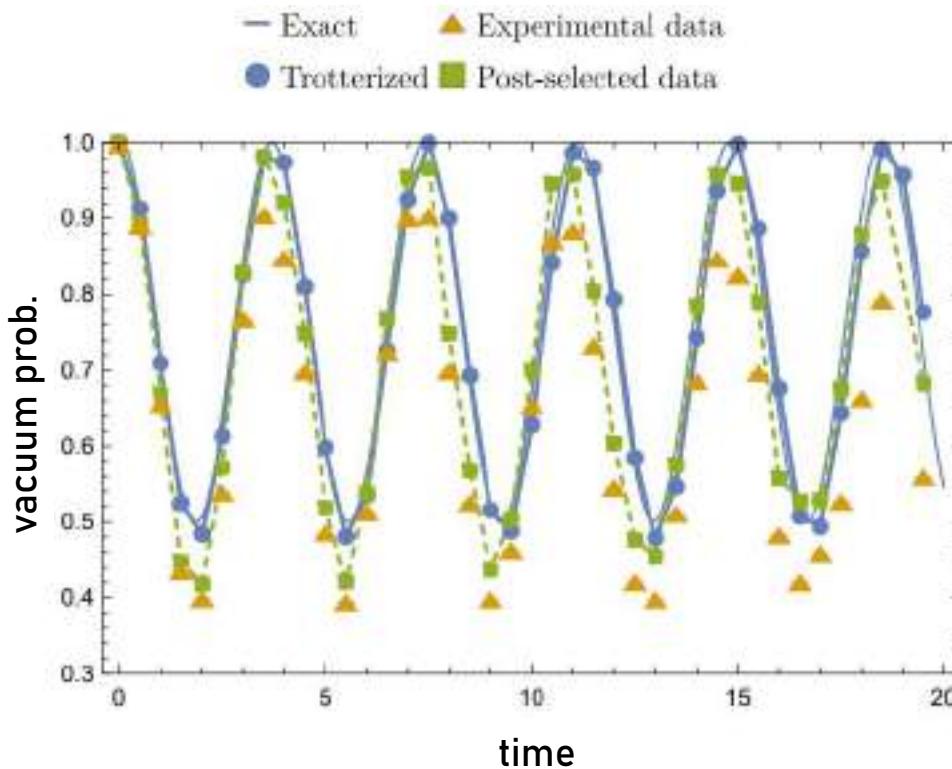
Digital Quantum Simulation: The Schwinger model

Results (1-site model, 2 qubits)



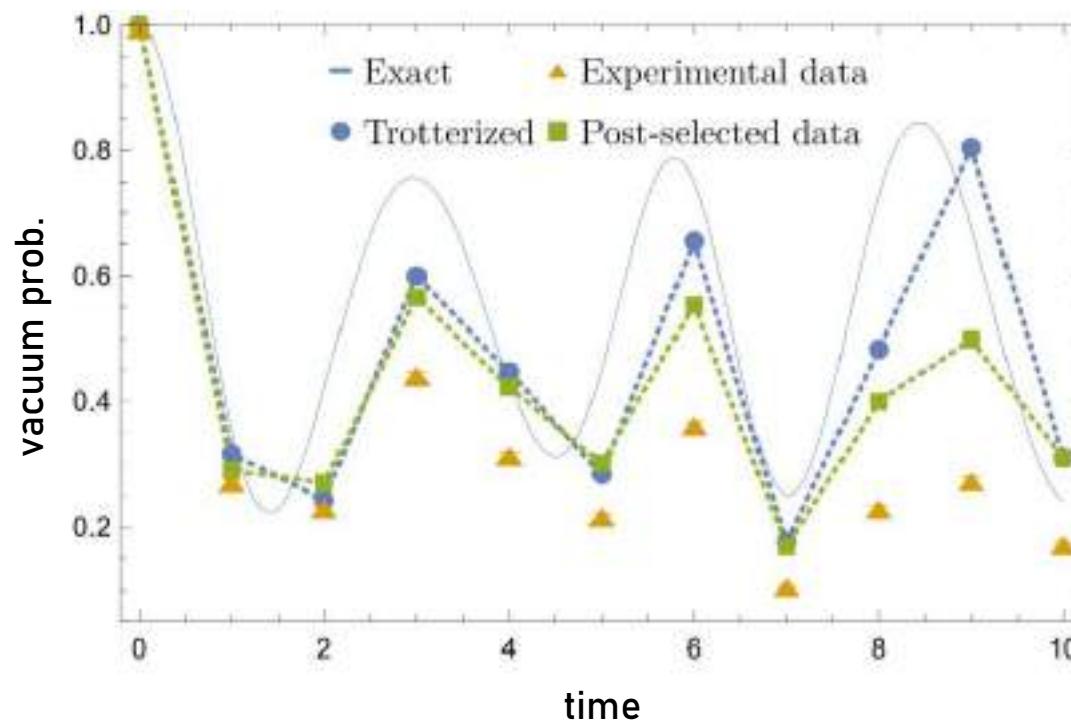
Digital Quantum Simulation: The Schwinger model

Results (1-site model, 2 qubits)



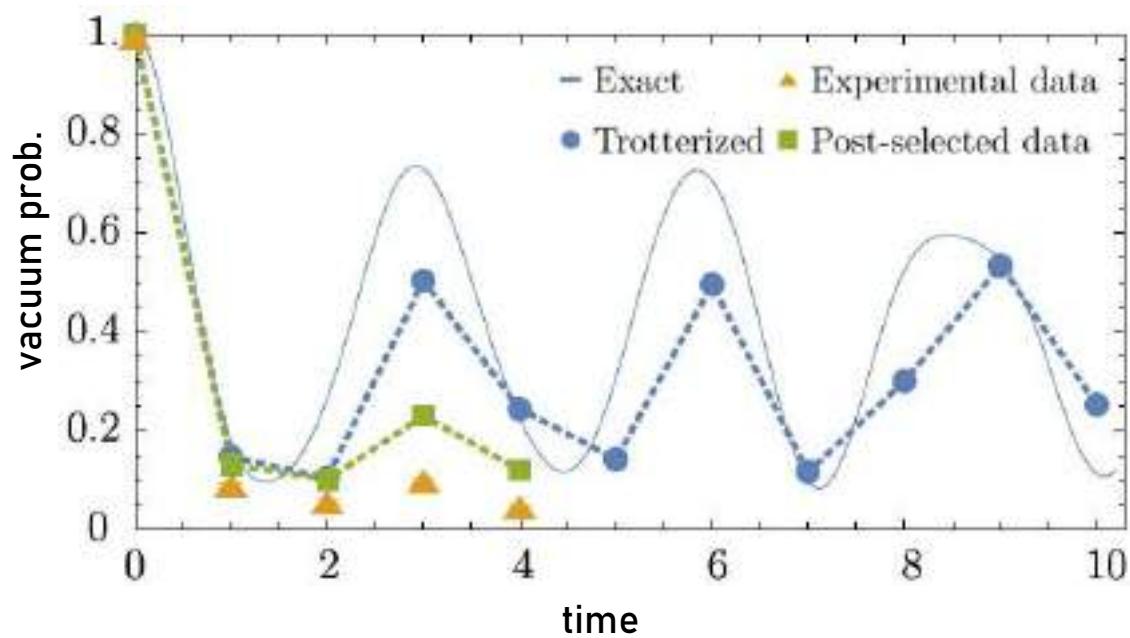
Digital Quantum Simulation: The Schwinger model

Results (2-site model, 4 qubits)



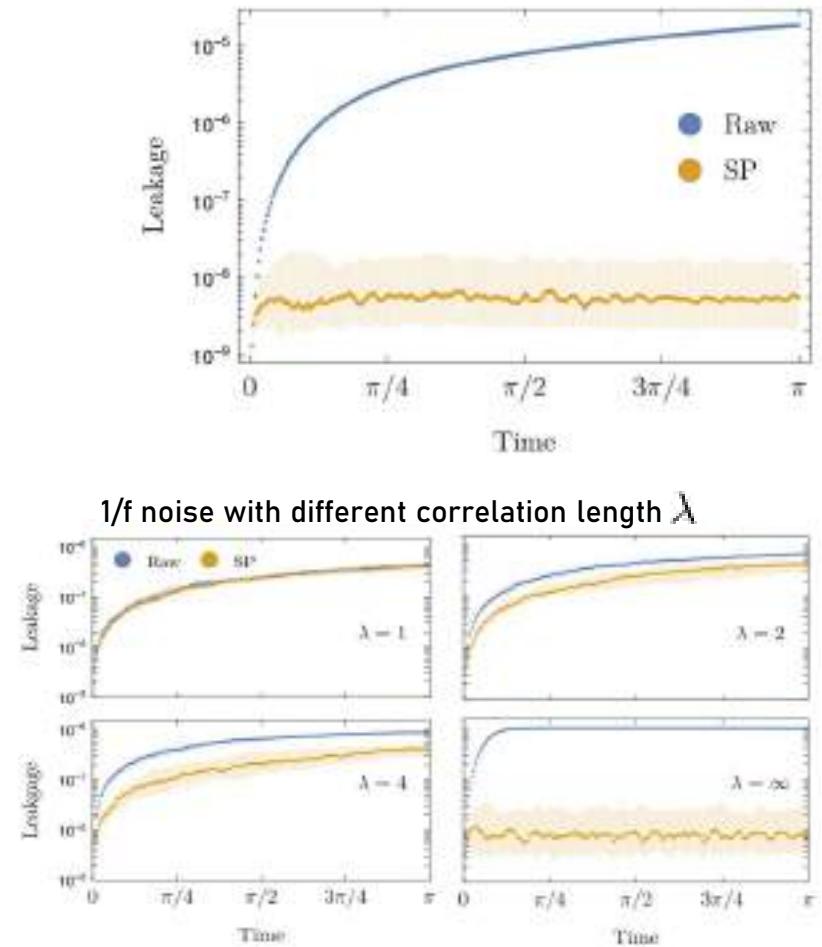
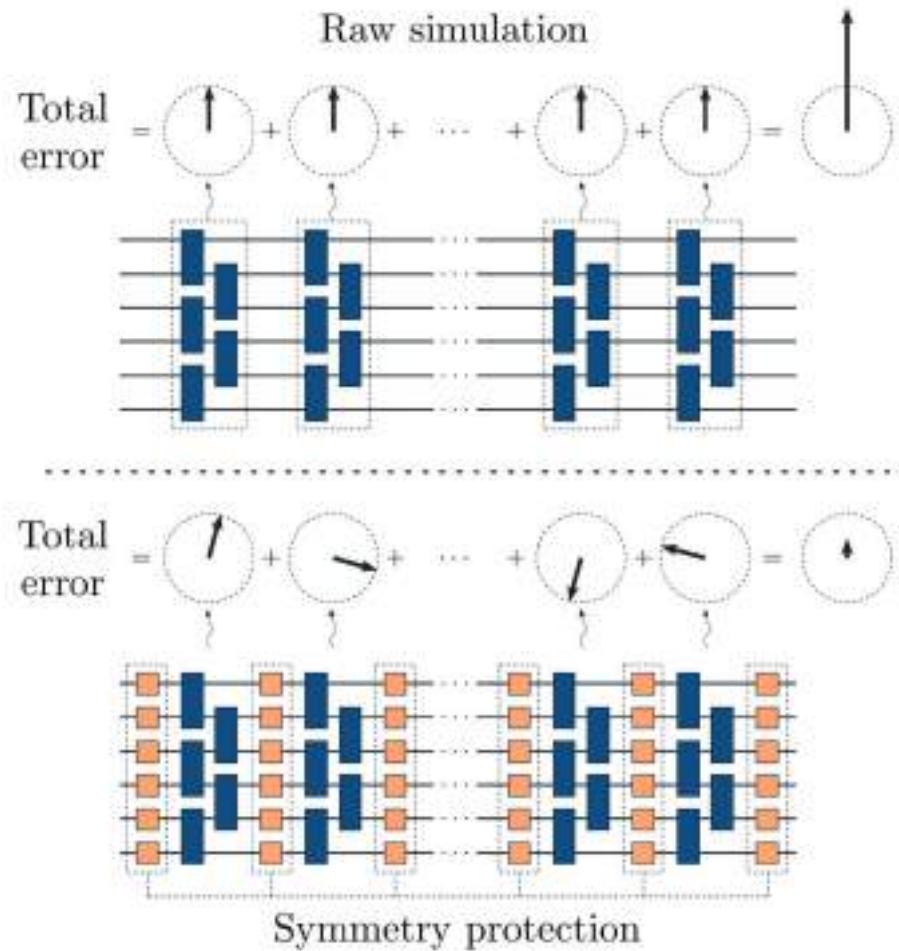
Digital Quantum Simulation: The Schwinger model

Results (3-site model, 6 qubits)

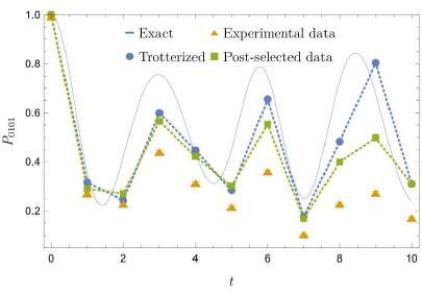
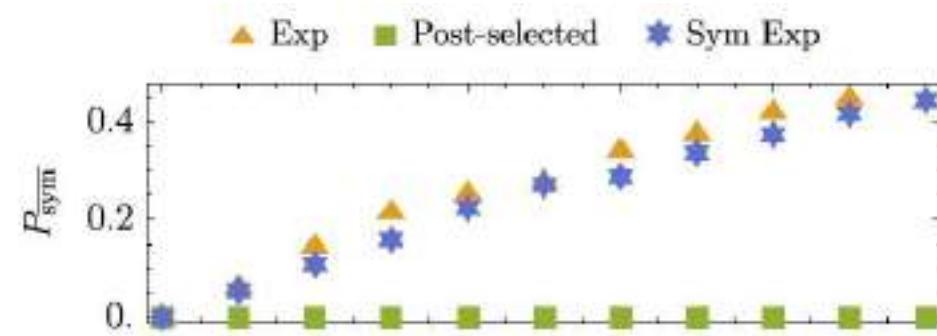


Digital Quantum Simulation: Symmetry Protection

Can we mitigate more of the errors?



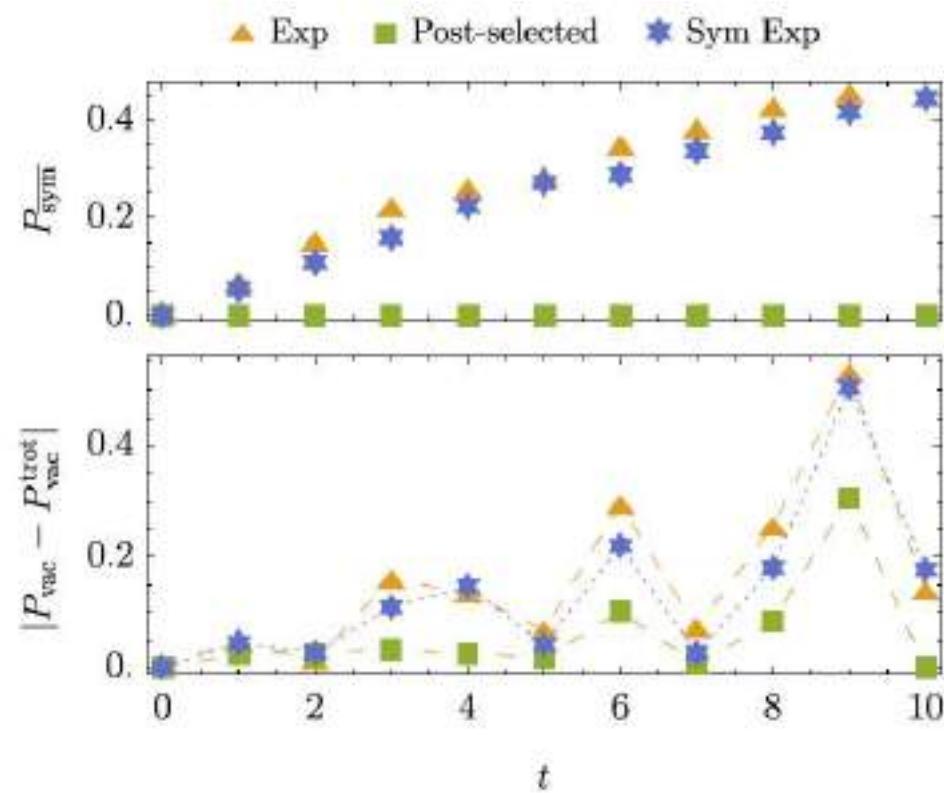
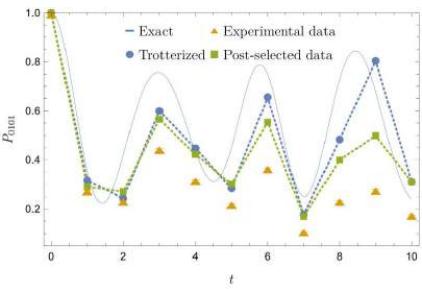
Digital Quantum Simulation: Symmetry Protection



Digital Quantum Simulation: Symmetry Protection

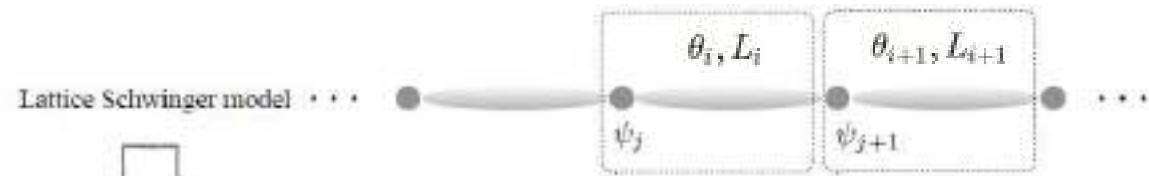
Conclusions:

- post selection is effective
- our errors is (time-)uncorrelated



Difficult Quantum Simulation... such as the Schwinger model

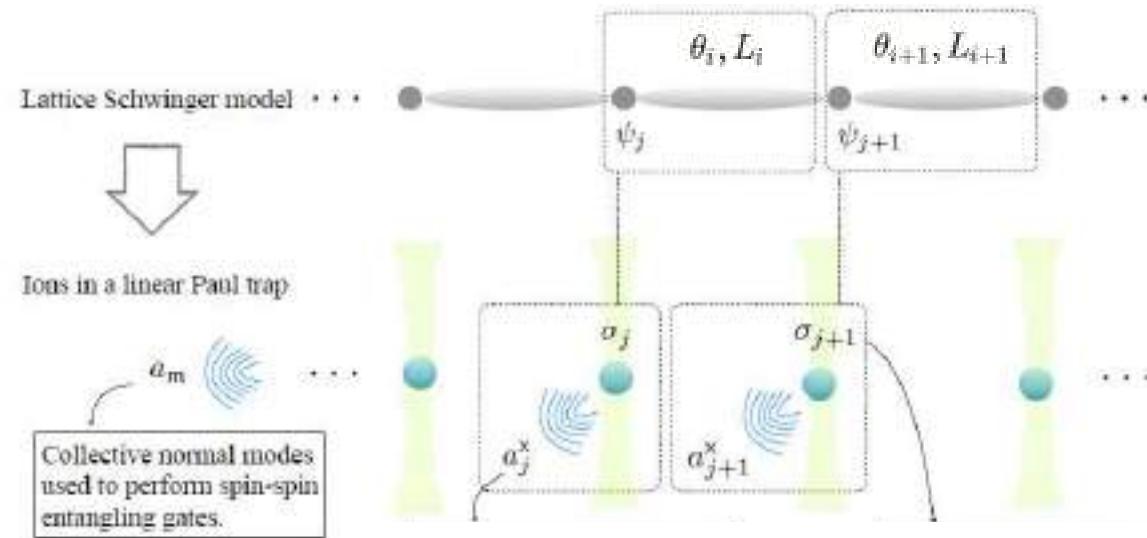
Analog Quantum Simulation Alternative ?



Zohreh Davoudi, NML, G. Pagano, Phys. Rev. Research 3, 043072 (2021)

Hybrid Quantum Simulation

The Schwinger model

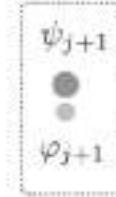
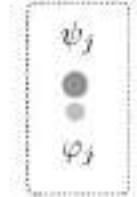


Zohreh Davoudi, NML, G. Pagano, Phys. Rev. Research 3, 043072 (2021)

Hybrid Quantum Simulation: The Yukawa model

$$V(r) = -\frac{g^2}{4\pi} \frac{1}{r} e^{-\mu r}$$

...



...

- Scalar field theory: scalar bosons (field) interact with fermions (matter)
- Describes the Higgs field interacting with the leptons and quarks
- Explains mass generation

$$H_{\text{Yukawa}} = H_{\text{Yukawa}}^{(I)} + H_{\text{Yukawa}}^{(II)} + H_{\text{Yukawa}}^{(III)}$$

$$H_{\text{Yukawa}}^{(III)} = gb \sum_{j=1}^N \psi_j^\dagger \varphi_j \psi_j$$

$$\varphi_j = \frac{1}{\sqrt{Nb}} \sum_{k=-N/2}^{N/2-1} \frac{1}{\sqrt{2\varepsilon_k}} (d_k^\dagger e^{-i2\pi kj/N} + d_k e^{i2\pi kj/N})$$

$$H_{\text{Yukawa}}^{(II)} = \sum_{k=-N/2}^{N/2-1} \varepsilon_k \left(d_k^\dagger d_k + \frac{1}{2} \right)$$

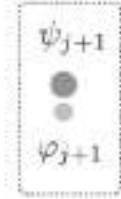
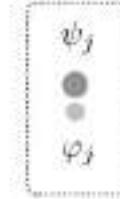
$$\varepsilon_k = \sqrt{\left(\frac{2\pi k}{Nb}\right)^2 + m_\varphi^2}$$

$$H_{\text{Yukawa}}^{(I)} = \sum_{j=1}^N \left[\frac{i}{2b} (\psi_j^\dagger \psi_{j+1} - \psi_{j+1}^\dagger \psi_j) + m_\psi (-1)^j \psi_j^\dagger \psi_j \right]$$

Hybrid Quantum Simulation: The Yukawa model

$$V(r) = -\frac{g^2}{4\pi} \frac{1}{r} e^{-\mu r}$$

...



...

- Scalar field theory: scalar bosons (field) interact with fermions (matter)
- Describes the Higgs field interacting with the leptons and quarks
- Explains mass generation

$$H_{\text{Yukawa}} = H_{\text{Yukawa}}^{(I)} + H_{\text{Yukawa}}^{(II)} + H_{\text{Yukawa}}^{(III)}$$

$$H_{\text{Yukawa}}^{(III)} = gb \sum_{j=1}^N \psi_j^\dagger \varphi_j \psi_j$$

Hybrid Quantum Simulation: The Yukawa model

$$V(r) = -\frac{g^2}{4\pi} \frac{1}{r} e^{-\mu r}$$



- Scalar field theory: scalar bosons (field) interact with fermions (matter)
- Describes the Higgs field interacting with the leptons and quarks
- Explains mass generation

$$H_{\text{Yukawa}} = H_{\text{Yukawa}}^{(I)} + H_{\text{Yukawa}}^{(II)} + H_{\text{Yukawa}}^{(III)}$$

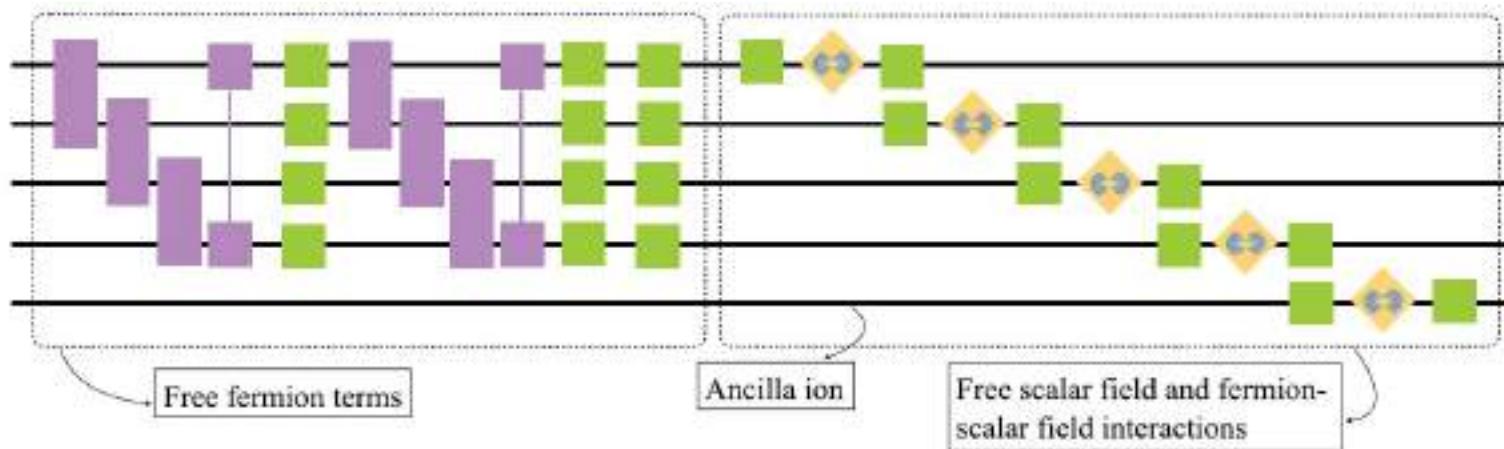
$$H_{\text{Yukawa}}^{(III)} = \sqrt{\frac{g^2 b}{8N}} \sum_{j=1}^N (\mathbb{I}_j + \sigma_j^z) \sum_{m=1}^N \underbrace{\frac{1}{\sqrt{\varepsilon_m}} (a_m^\dagger e^{-i\frac{2\pi j}{N}(m-\frac{N}{2}-1)} + a_m e^{i\frac{2\pi j}{N}(m-\frac{N}{2}-1)})}_{\begin{array}{l} \text{ancilla ion} \\ \text{red & blue sideband (rwa, detuned)} \end{array}}$$

Hybrid Quantum Simulation: The Yukawa model

Nhung Nguyen



Map fermions to spins and bosons to phonons



Single-spin gates



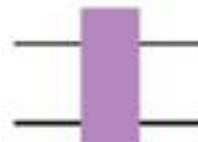
$$\{R_j^{\sigma}(\theta_j, \phi_j), R_j^{\pm}(\theta_j)\}$$

Spin-(normal)
phonon gate

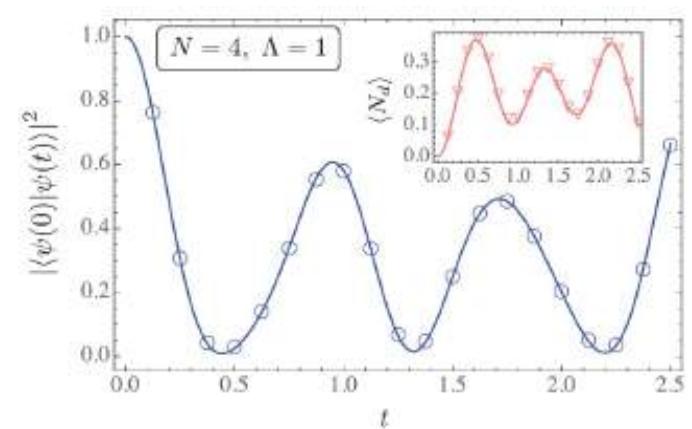


$$R_{k,j}^{\sigma\alpha}(\{\theta_{k,j}\}, \{\phi_{k,j}\})$$

Two-spin gate (MS)



$$R_{j,j'}^{\sigma\sigma}(\theta_{j,j'})$$

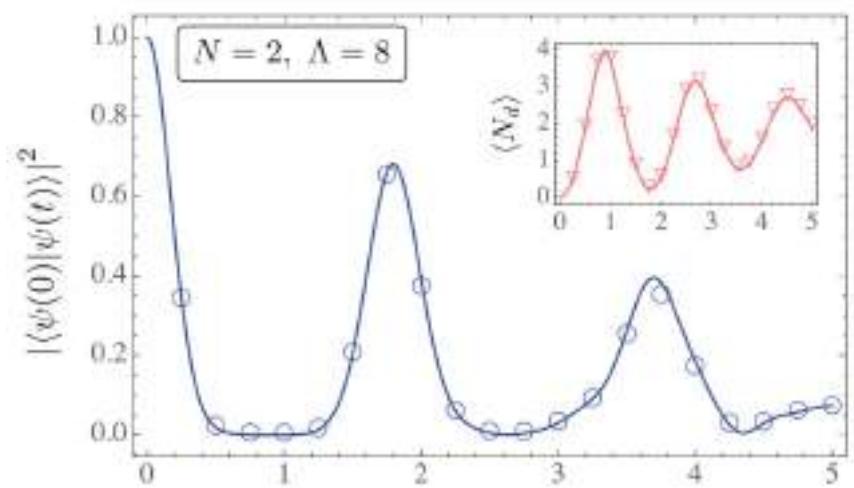
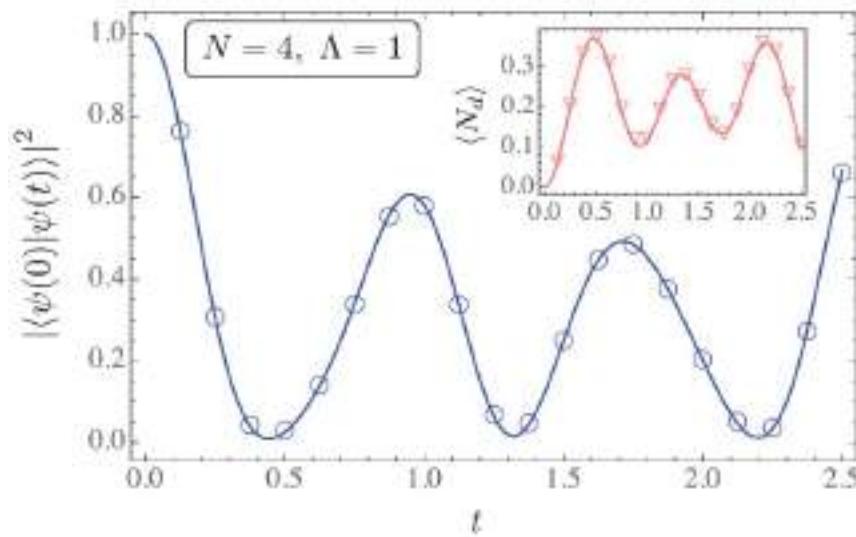


Hybrid Quantum Simulation: The Yukawa model

$\psi(0)$

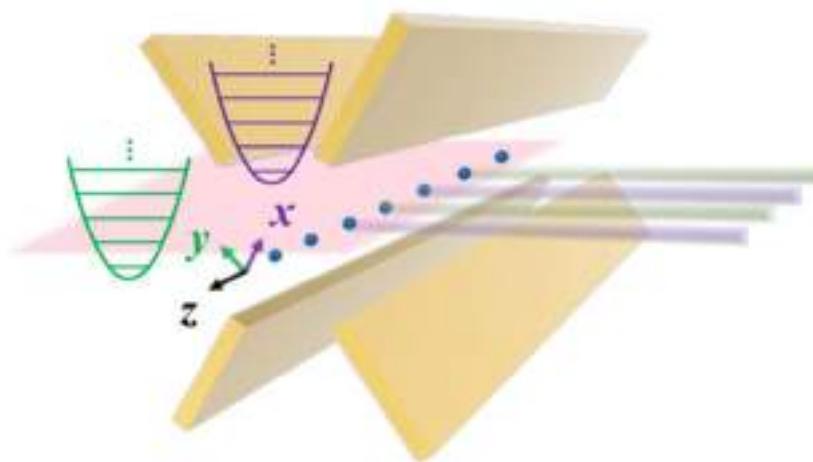
- Start in the ground state of non-interacting H (no fermions, anti-fermions, or bosons)
- Time evolution with Trotterization
- Measure revivals, avg. boson number

$$|\langle \psi(0) | \psi(t) \rangle|^2 \quad \langle N_d \rangle \equiv \frac{1}{N} \langle \psi(t) | \sum_{k=-N/2}^{N/2-1} d_k^\dagger d_k | \psi(t) \rangle$$

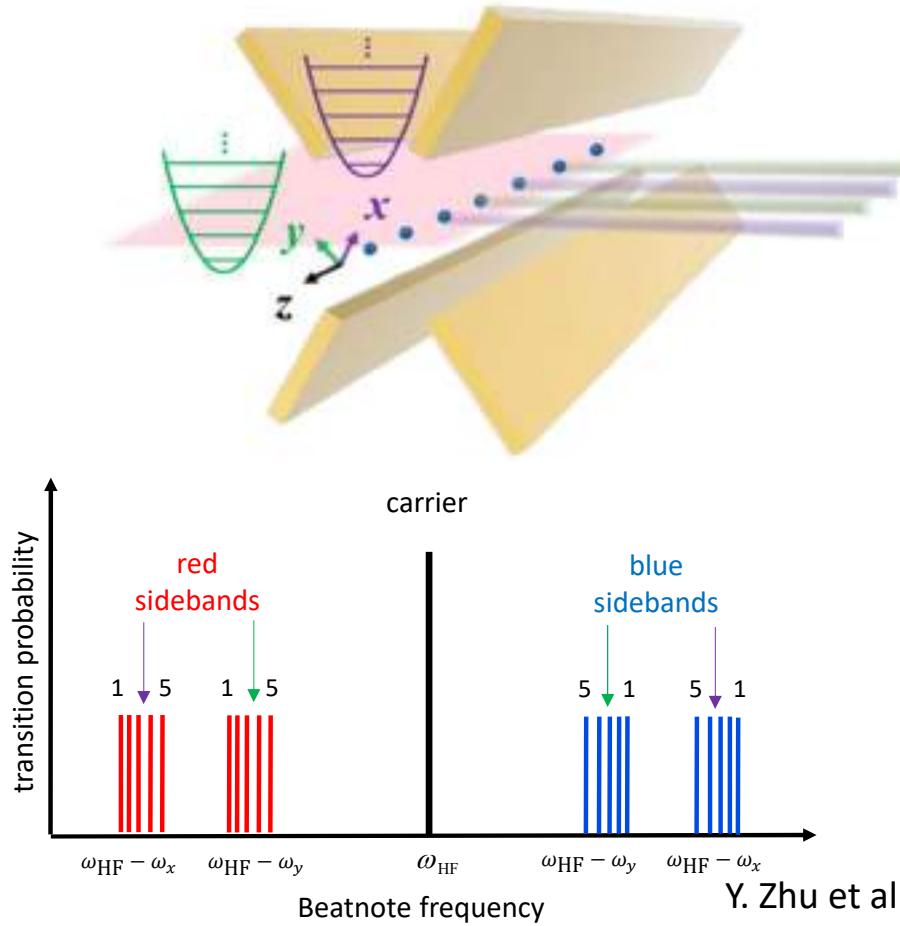


Pairwise parallel gates

Pairwise parallel gates

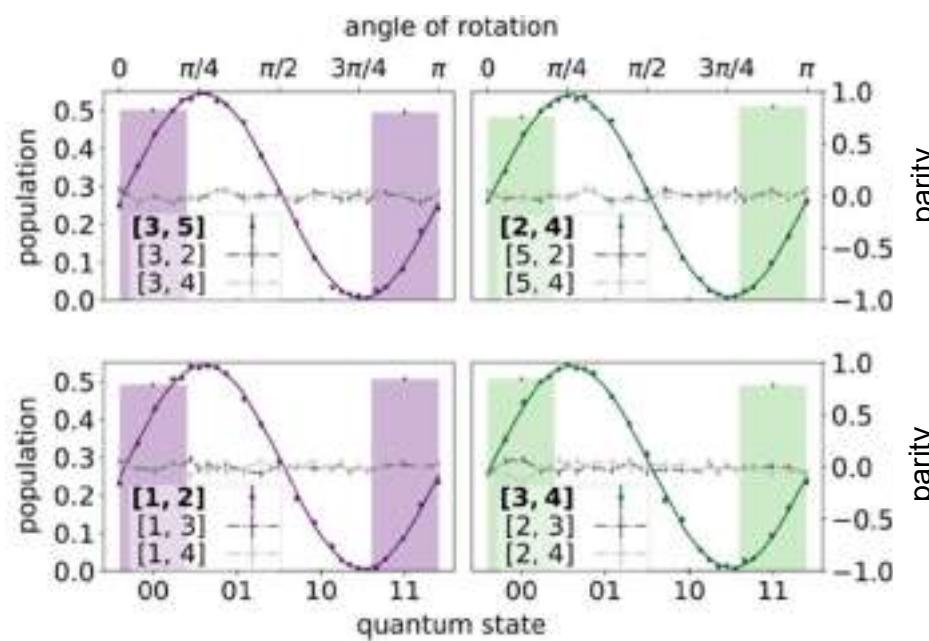


Pairwise parallel gates



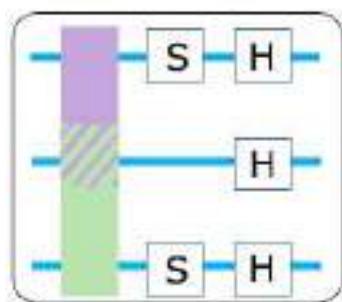
Pairwise parallel gates

$$|00\rangle \rightarrow \frac{1}{\sqrt{2}}(|00\rangle - i|11\rangle)$$



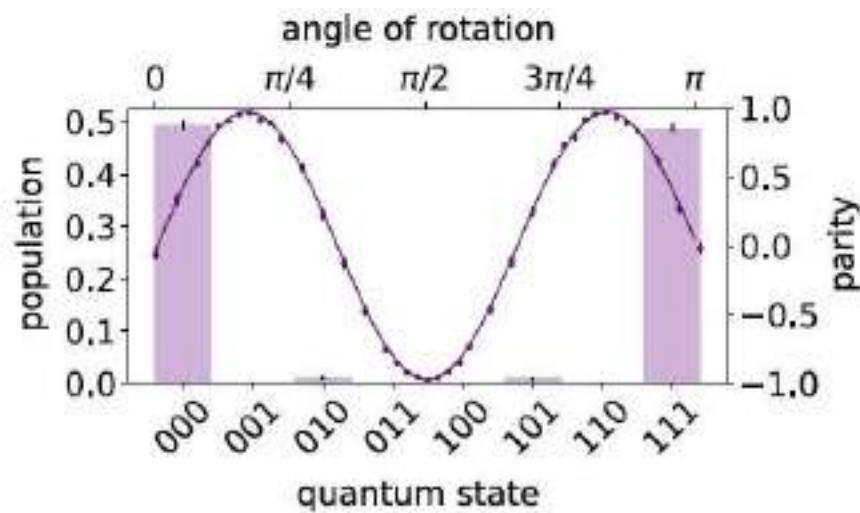
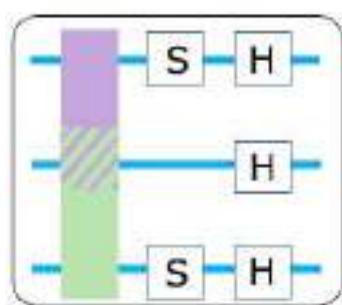
Pairwise parallel gates

Parallel gates with overlapping ion: GHZ state



Pairwise parallel gates

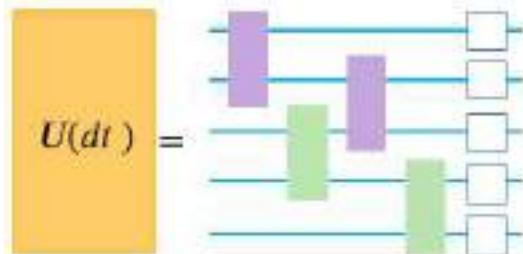
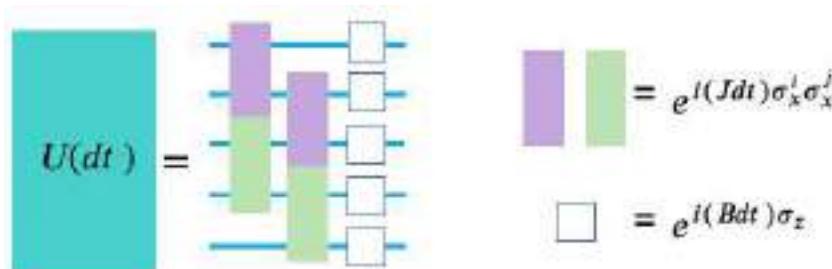
Parallel gates with overlapping ion: GHZ state



Pairwise parallel gates

Transverse-field Ising model:

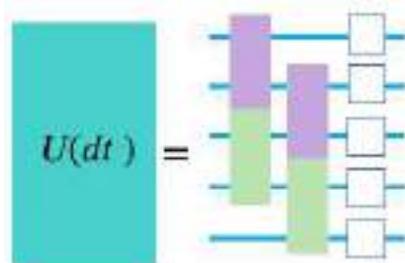
$$H = -J \sum_{i=1}^4 \sigma_x^i \sigma_x^{i+1} - B \sum_{i=1}^5 \sigma_z^i$$



Pairwise parallel gates

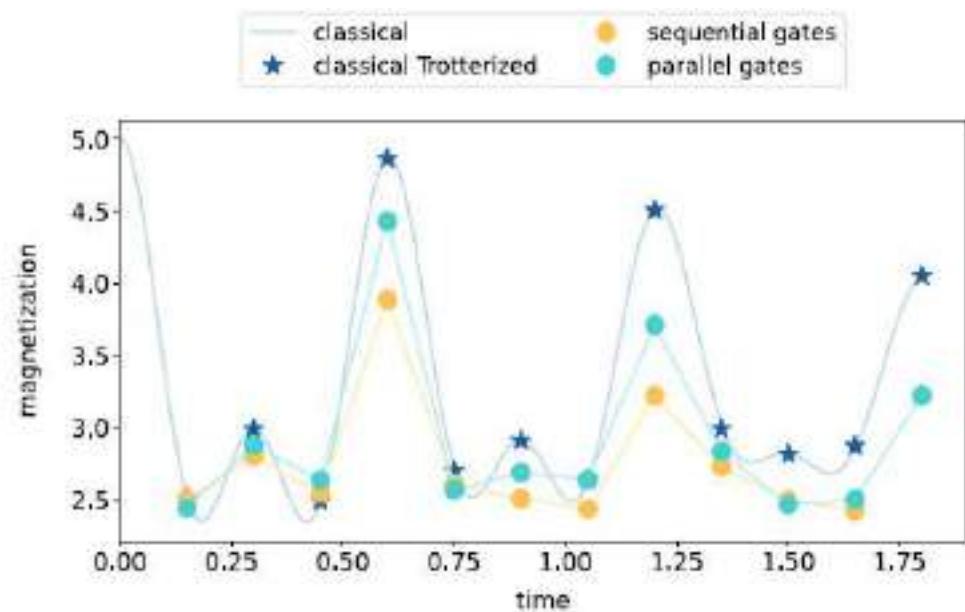
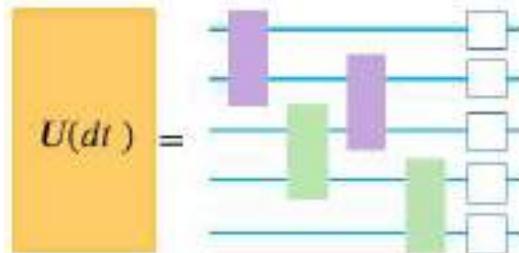
Transverse-field Ising model:

$$H = -J \sum_{i=1}^4 \sigma_x^i \sigma_x^{i+1} - B \sum_{i=1}^5 \sigma_z^i$$



$$\begin{array}{c} \text{purple block} \\ \text{green block} \end{array} = e^{i(Jdt)\sigma_x^i\sigma_x^j}$$

$$\square = e^{i(Bdt)\sigma_z}$$



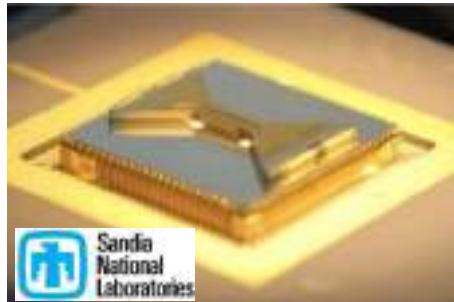
Outlook: A new ion trap platform

collaboration with G. Pagano (Rice)

Outlook: A new ion trap platform

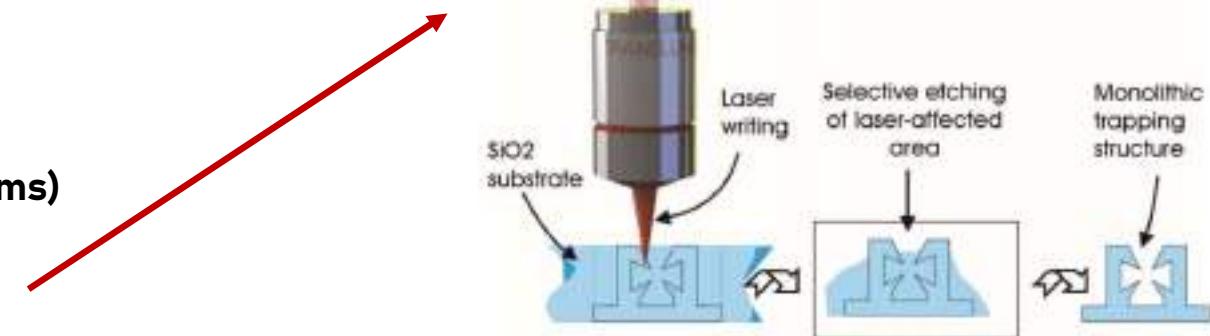


3D blade trap (with problems)



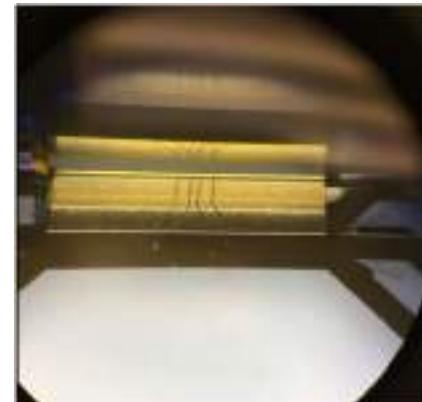
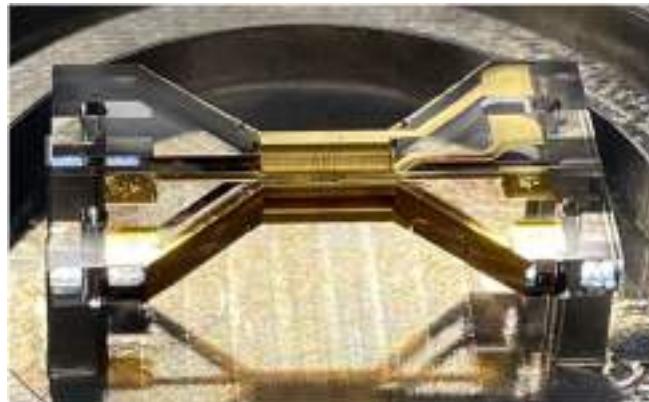
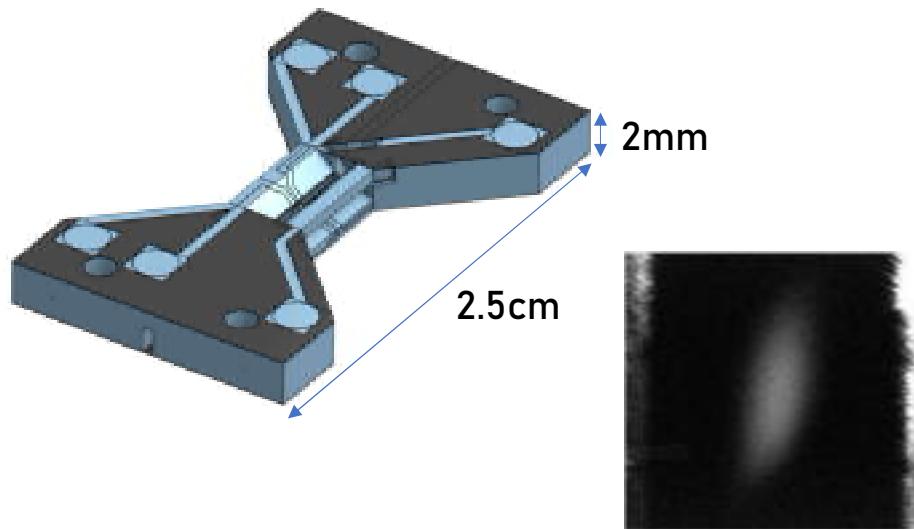
2D surface traps (with fewer/different problems):
DQC groups, PTB, NICT, Honeywell, NIST etc.

Monolithic 3D trap made of Fused Silica by Translume Inc.



collaboration with G. Pagano (Rice)

Outlook: A new ion trap platform



TAMOS to be est. Oct. 2023



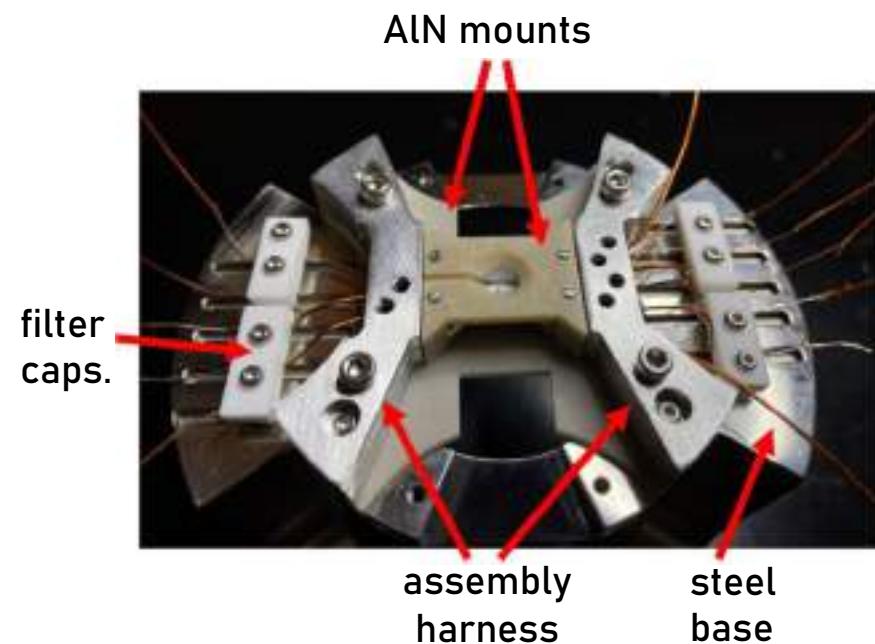
R. Zhuravel



G. Pagano



NML





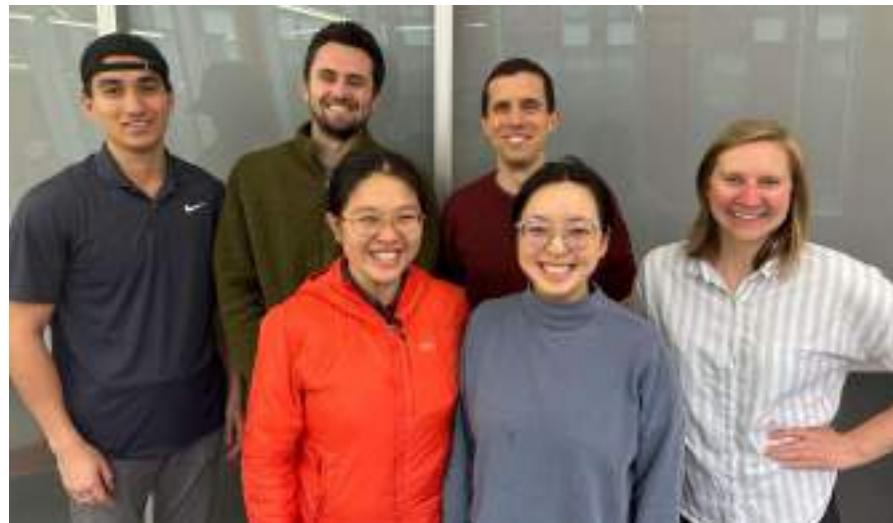
Tuesday poster 19

Elijah
Mossman (u)

Liam
Jeanette

NML

Alaina
Green



Nhung
Nguyen

Yingyue
Zhu

Monday poster 27

Denton
Wu
Minseo
Kim (u)
Devon
Valdez
Henry
Luo
Yuanheng
Xie
Jiqing
Fan (u)
NML



Michael
Straus

Mika
Chmielewski

Xinyi
Dai

Sophie
Decoppet

Ana
Ferrari

Monday poster 18



QUANTUM SYSTEMS ACCELERATOR
Catalyzing the Quantum Ecosystem



Institute for
Robust Quantum
Simulation





Duke Quantum Center

Monday posters: 152 (M. Donofrio), 163 (S. Patel), 177 (K. Ranawat),
147 (E. Reed), 176 (J. Whitlow)

Tuesday posters: 60 (J. O'Reilly), 209 (D. Biswas), 113 (J. Toast)



Brown



Calderbank



Cetina



Kim



Kozhanov



Linke



Marvian



Monroe



Noel